

# Excited radiative transitions in charmonium from lattice QCD

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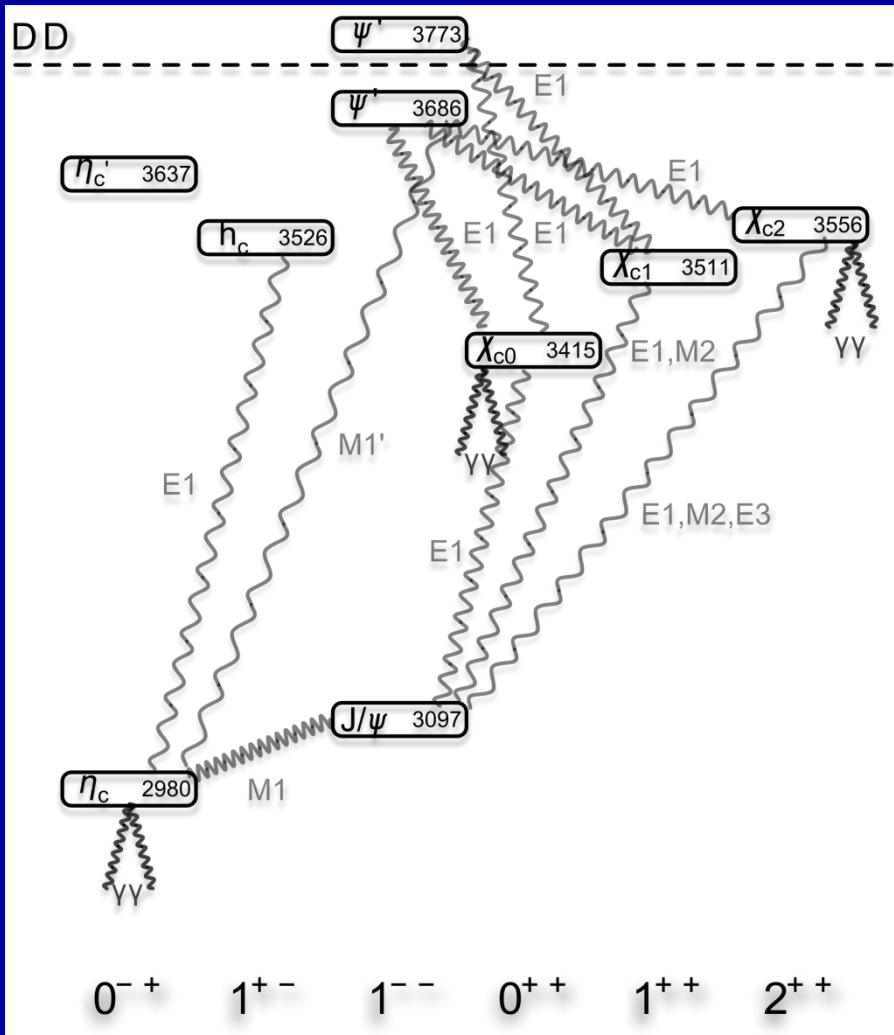
QWG Workshop, Fermilab, May 2010

With Jo Dudek and Robert Edwards  
Previous work: Nilmani Mathur, David Richards and  
Ermal Rrapaj (and *Hadron Spectrum Collaboration*)

# Outline

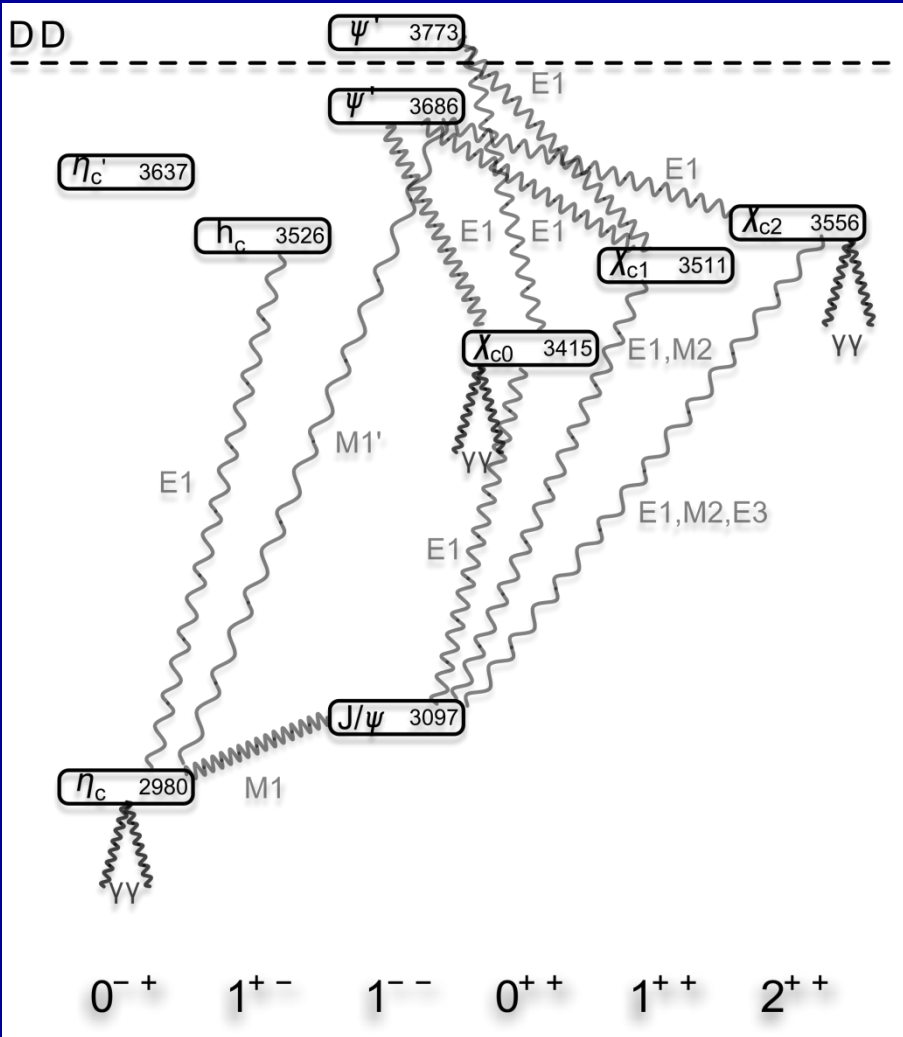
- Introduction and motivation
- Method
- Result highlights and interpretations
- Summary and outlook

# Charmonium radiative transitions



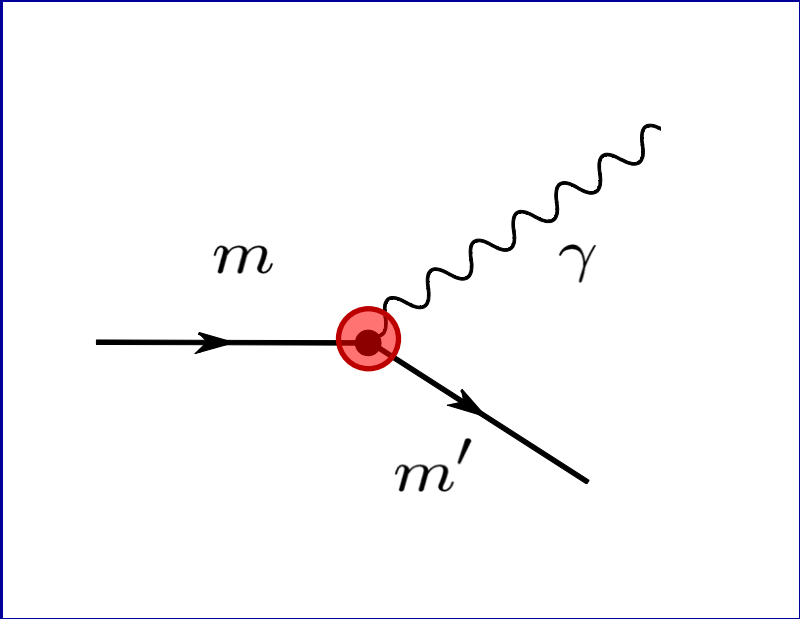
Below DD threshold radiative transitions have significant BRs

# Charmonium radiative transitions

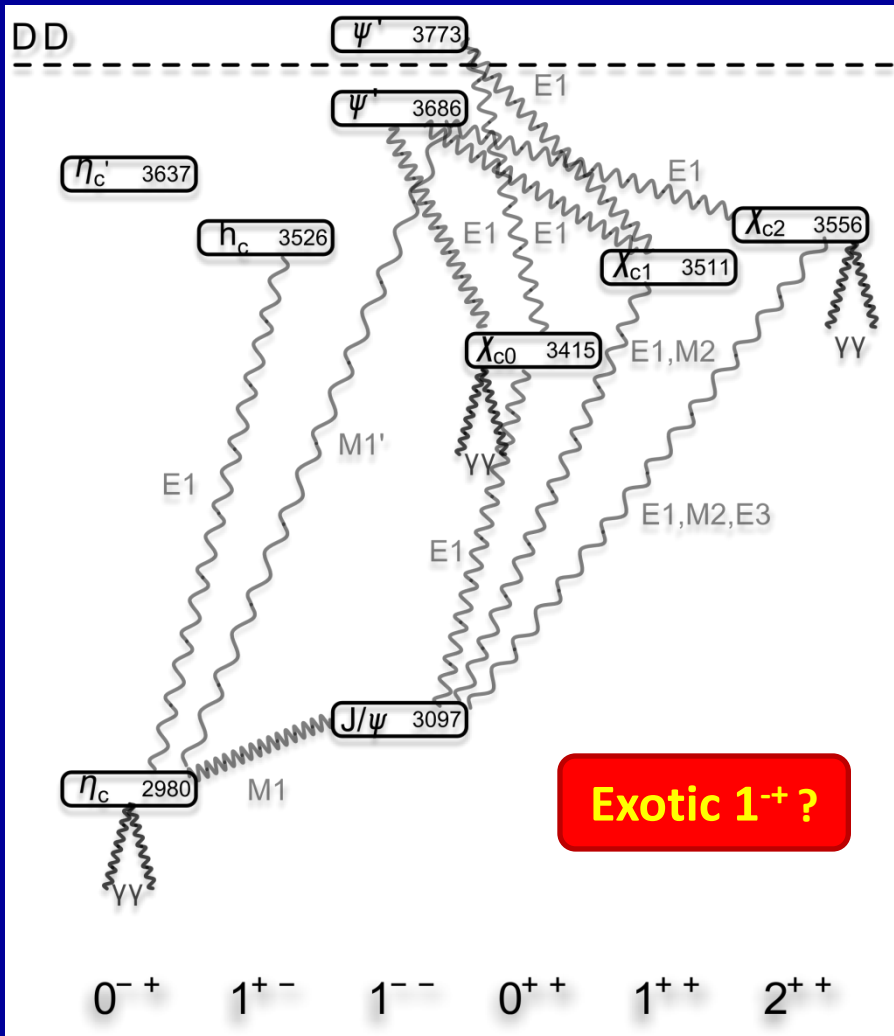


Below DD threshold radiative transitions have significant BRs

## Meson – Photon coupling

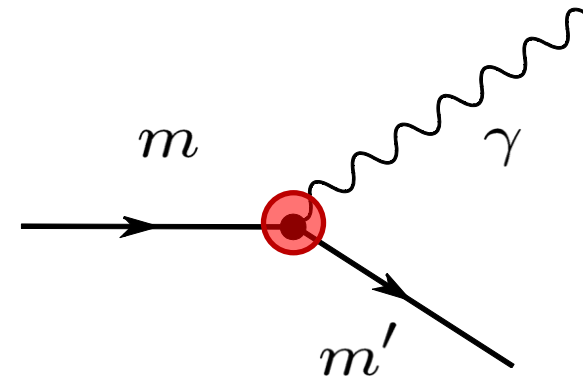


# Charmonium radiative transitions



Below DD threshold radiative transitions have significant BRs

Meson – Photon coupling



# Broader Picture

Develop Lattice QCD techniques

Test in the charmonium system

Apply to lighter mesons...

# Broader Picture

Develop Lattice QCD techniques

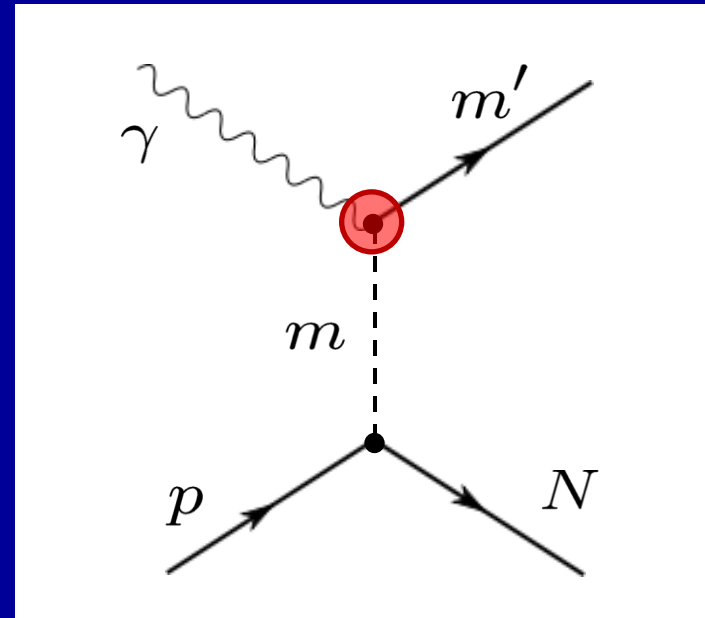
Test in the charmonium system

Apply to lighter mesons...

Photoproduction at GlueX  
(JLab 12 GeV upgrade)

Exotic  $1^{-+}$  ?

Progress on the light meson spectrum in Dudek et al  
PRL103 262001 (2009) and arXiv:1004.4930.



# Photocouplings on the lattice

Two-point correlation functions with a large basis of operators  $\rightarrow$   
energies and matrix elements (Z)

$$O(t) = \sum_{\vec{x}} e^{i\vec{p} \cdot \vec{x}} \bar{\psi}(x) \Gamma_i \overleftrightarrow{D}_j \overleftrightarrow{D}_k \dots \psi(x)$$

$$C(t) = \langle 0 | O_i(t) O_j(0) | 0 \rangle$$

$$Z_i^{(n)} \equiv \langle 0 | O_i | n \rangle$$



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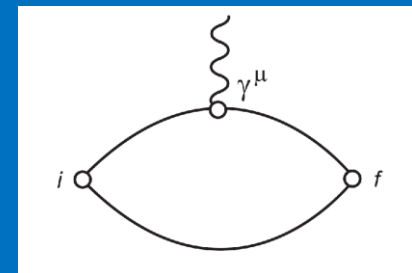
$$C(t) = \langle 0 | O_i(t) O_j(0) | 0 \rangle$$

$$Z_i^{(n)} \equiv \langle 0 | O_i | n \rangle$$

$$C_{ij}(t_f, t, t_i) = \langle 0 | O_i(t_f) \bar{\psi}(t) \gamma^\mu \psi(t) O_j(t_i) | 0 \rangle$$

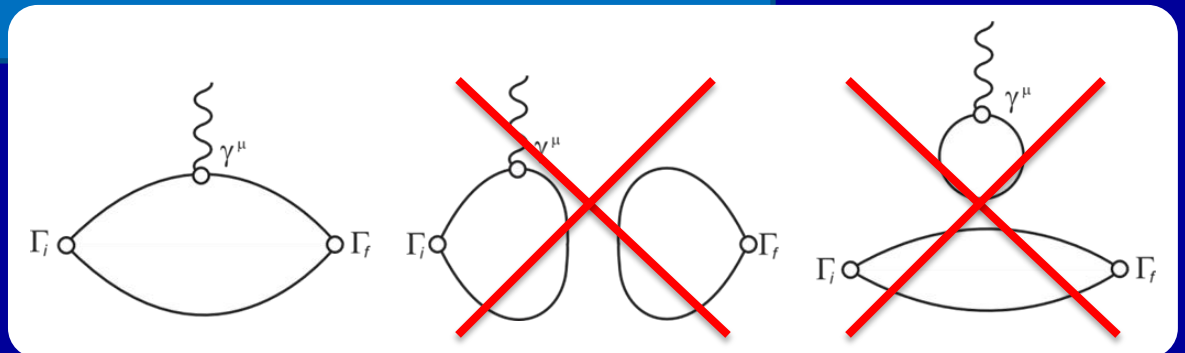
Photocouplings from three-point correlators

Need energies and Z's from two-point analysis



# Charmonium radiative transitions

- Caveats:
  - Quenched (no quark loops; no light quarks at all)
  - One lattice spacing ( $a_t^{-1} = 6.05$  GeV)
  - One volume ( $L_s \approx 1.2$  fm)
  - Only connected diagrams



Only some highlights here; more results and details in  
Dudek, Edwards & CT, PR **D79** 094504 (2009)

Also: Dudek et al PR **D77** 034501 (2008); Dudek & Rrapaj PR **D78** 094504 (2008)

# Exotic $1^{-+}$

Spectrum analysis:  $1^{-+} \eta_{c1}$  state found at 4300(50) MeV

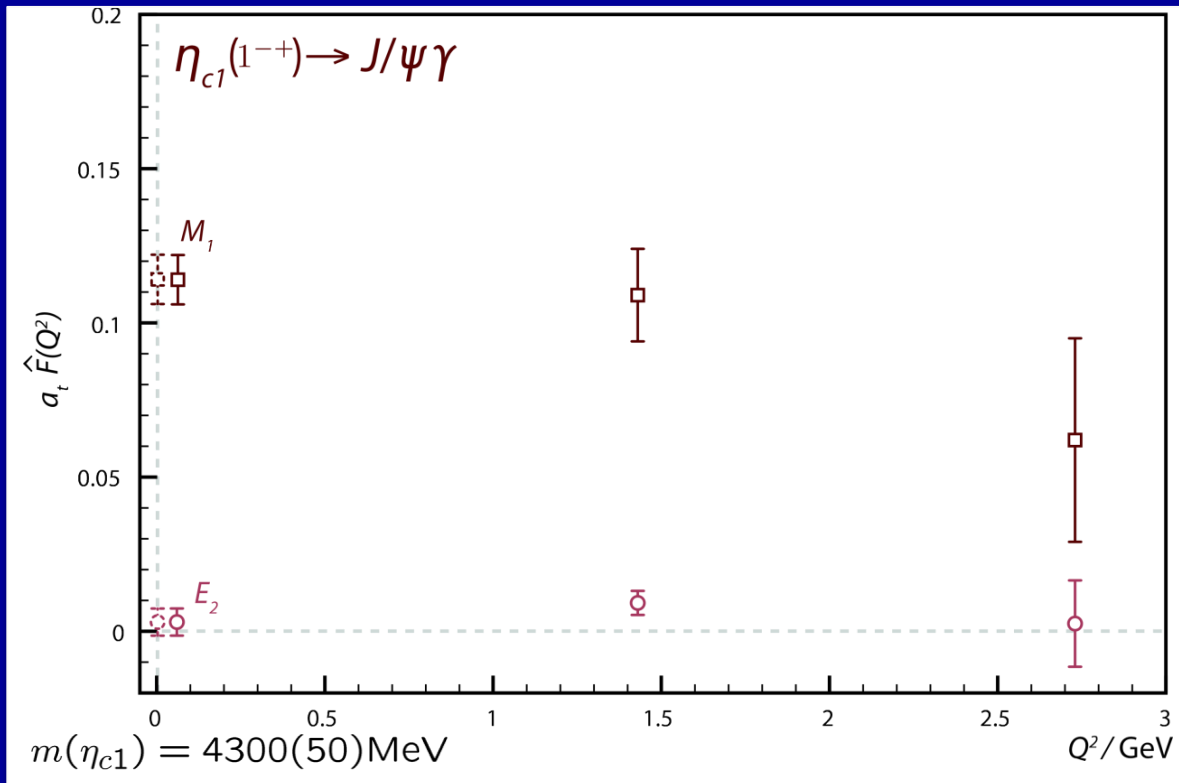
Exotic quantum numbers – can't be fermion-antifermion pair

Can't be a molecular/multi-quark state in **quenched** lattice calc.

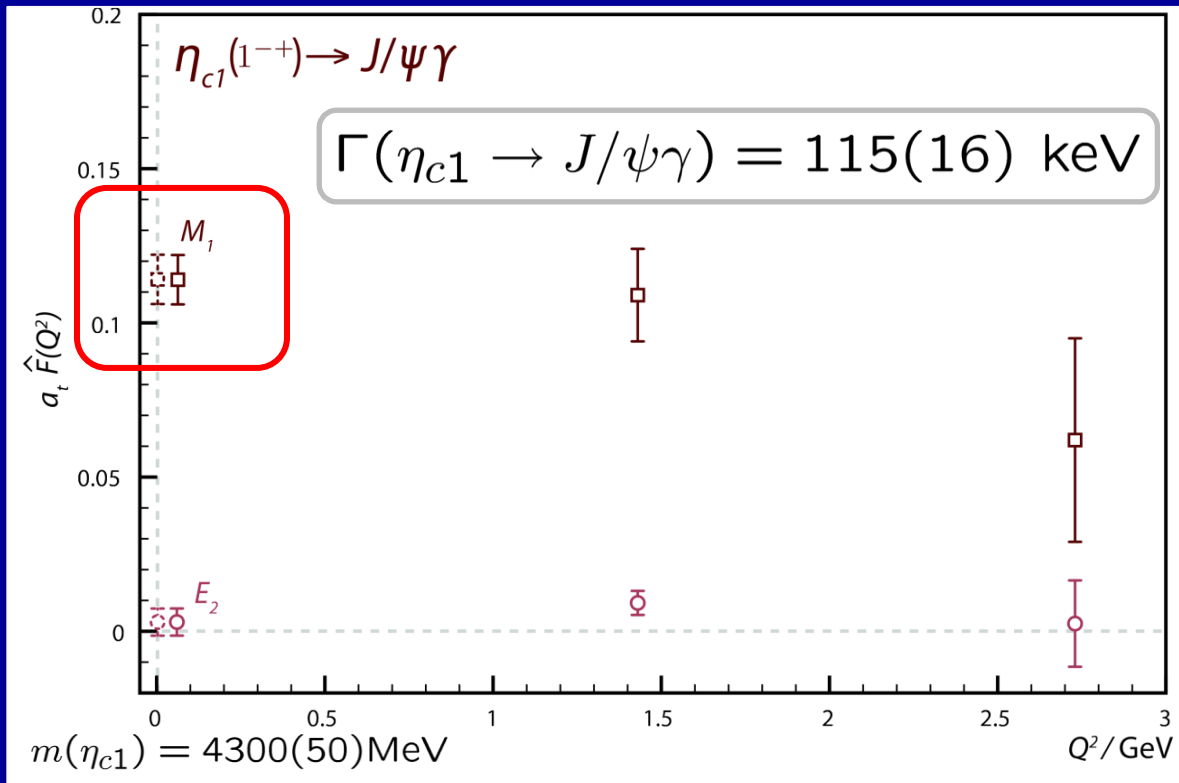
→ Strongly suggests a hybrid

What about radiative transitions?

# Exotic $1^{--}$ – Vector $1^{--}$



# Exotic $1^{--}$ – Vector $1^{--}$



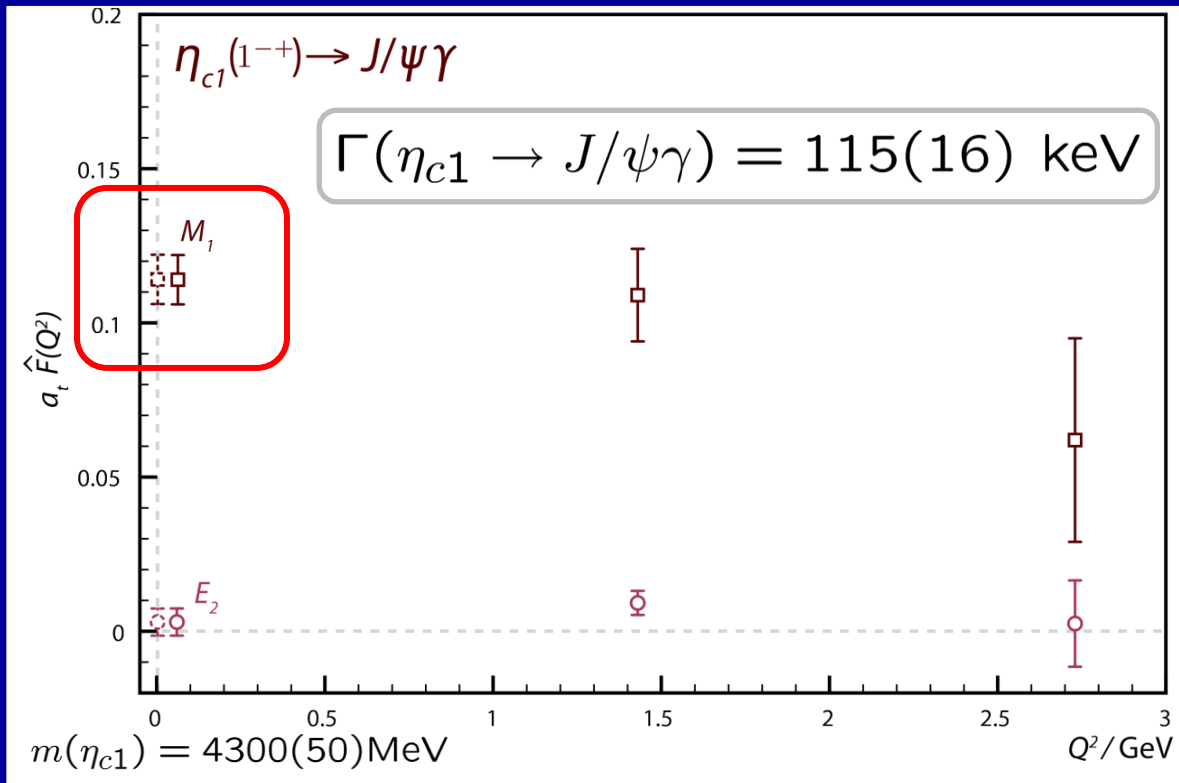
$M_1$  multipole dominates

Same scale as many measured conventional charmonium transitions

BUT very large for an  $M_1$  transition

$\Gamma(J/\psi \rightarrow \eta_c \gamma) \sim 2 \text{ keV}$

# Exotic $1^{--}$ – Vector $1^{--}$



$M_1$  multipole dominates

Same scale as many measured conventional charmonium transitions

BUT very large for an  $M_1$  transition

$\Gamma(J/\psi \rightarrow \eta_c \gamma) \sim 2 \text{ keV}$

- Usually  $M_1 \rightarrow$  spin flip (e.g.  $^3S_1 \rightarrow ^1S_0$ )  $\rightarrow 1/m_c$  suppression
- Spin-triplet hybrid  $\rightarrow$  extra gluonic degrees of freedom  
 $\rightarrow M_1$  transition without spin flip  $\rightarrow$  not suppressed

# Tensor $2^{++}$ – Vector $1^-$

$E_1, M_2, E_3$

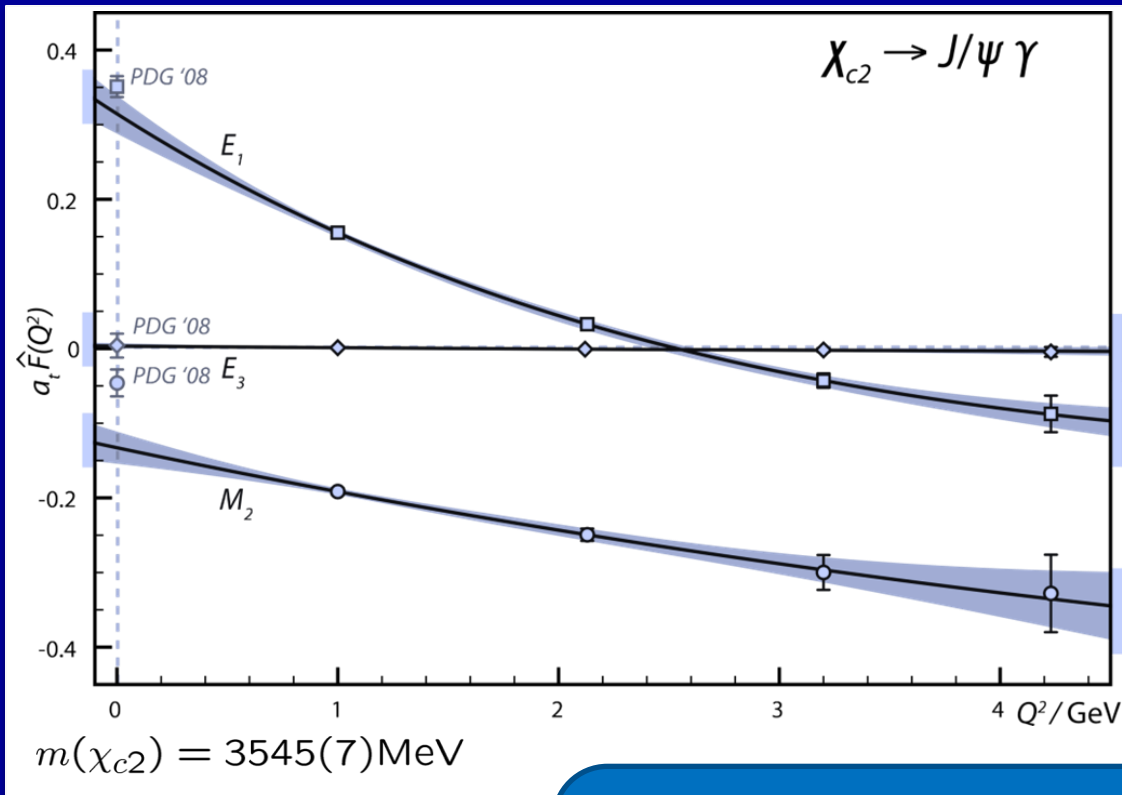
Three tensor states found in spectrum analysis:

$$\chi_{c2} \approx 3545(7) \text{ MeV}$$

$$\chi'_{c2} \approx 4115(28) \text{ MeV}$$

$$\chi''_{c2} \approx 4165(30) \text{ MeV}$$

# Tensor $2^{++}$ – Vector $1^-$



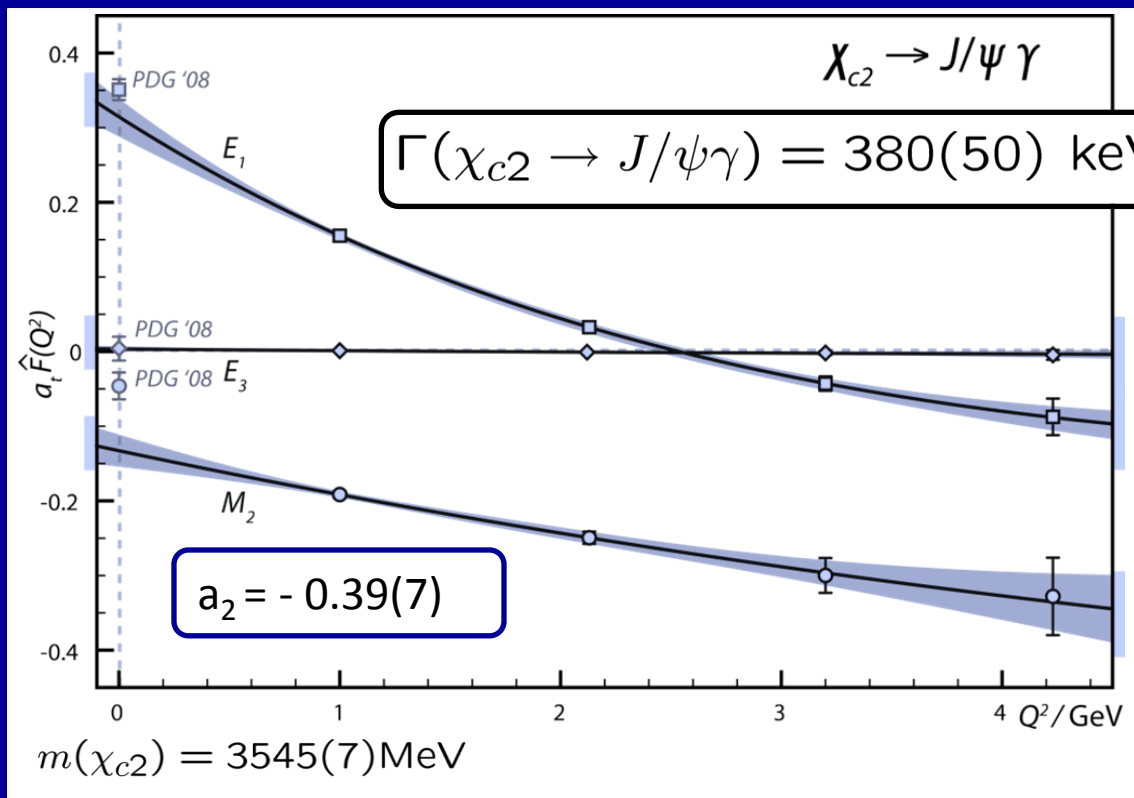
$E_1, M_2, E_3$

- Lattice: discrete set of allowed momenta
- Can't calculate at  $Q^2 = 0$  and so extrapolate:

$$F_k(Q^2) = F_k(0) \left(1 + \lambda Q^2\right) e^{-\frac{Q^2}{16\beta^2}}$$



# Tensor $2^{++}$ – Vector $1^-$

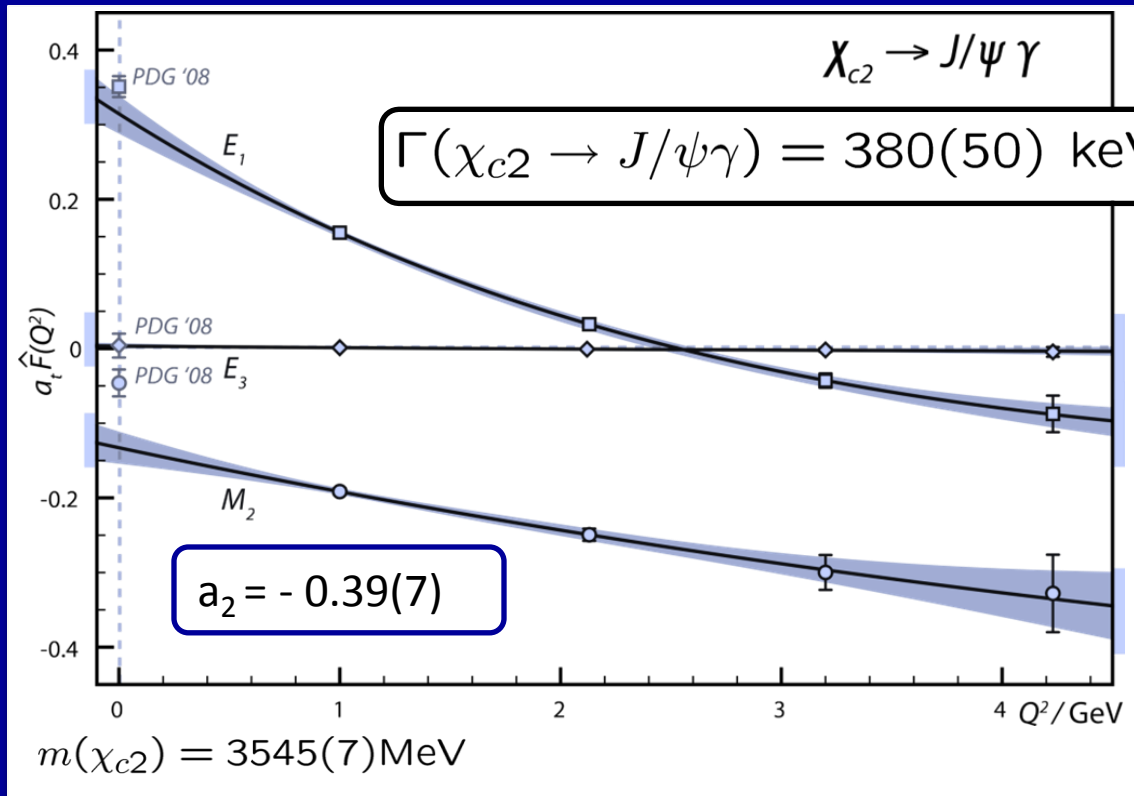


$E_1, M_2, E_3$

PDG08: 406(31) keV

Quark models ( $1^3P_2$ )  
~ 290 – 420 keV

# Tensor $2^{++}$ – Vector $1^-$



$E_1, M_2, E_3$

PDG08: 406(31) keV

Quark models ( $1^3P_2$ )  
~ 290 – 420 keV

$$a_2 = M_2 / \sqrt{(E_1^2 + M_2^2 + E_3^2)}$$

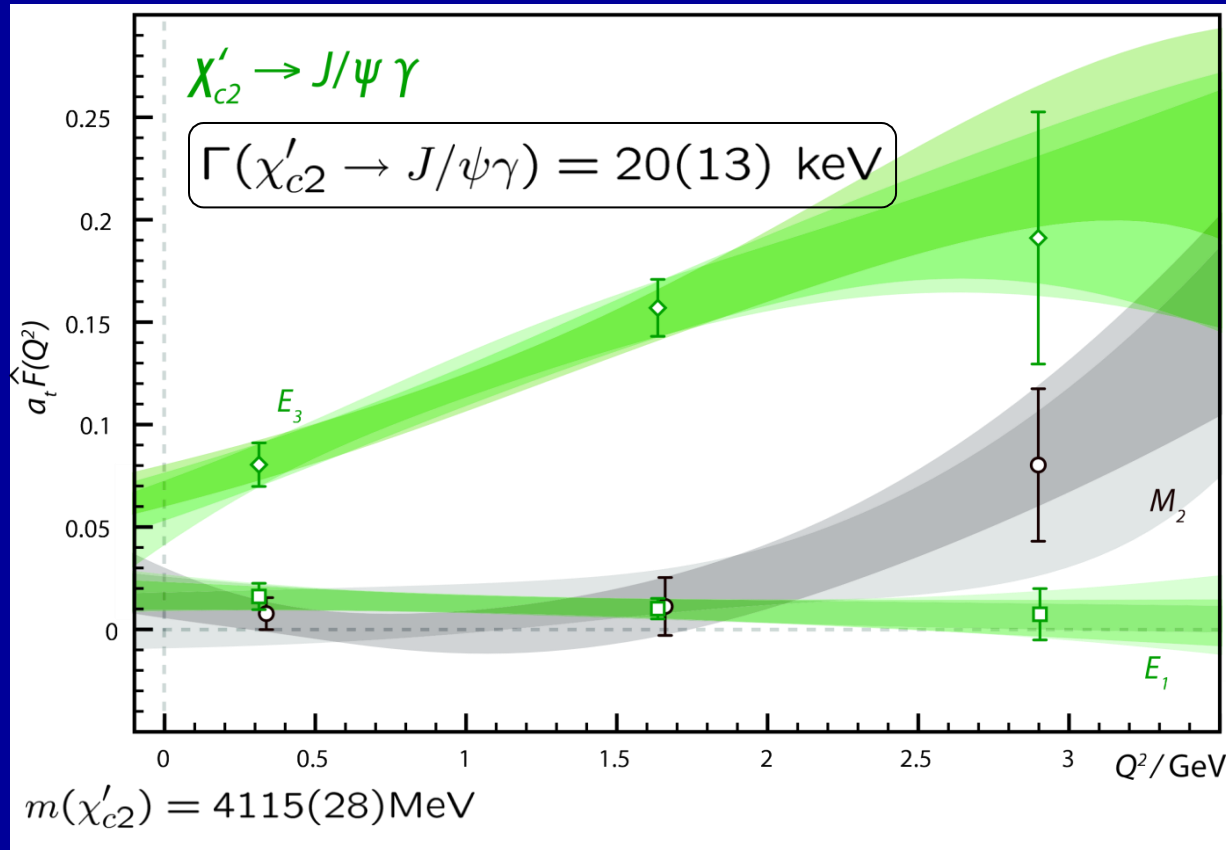
PDG08: -0.13(5)

CLEO: -0.079(19)

[CLEO PRD80 112003 (2009)]

- Same hierarchy as expected:  $|E_1(0)| > |M_2(0)| \gg |E_3(0)|$
- Ratio  $|M_2/E_1|$  is considerably larger than experiment

# Tensor $2^{++}$ – Vector $1^-$

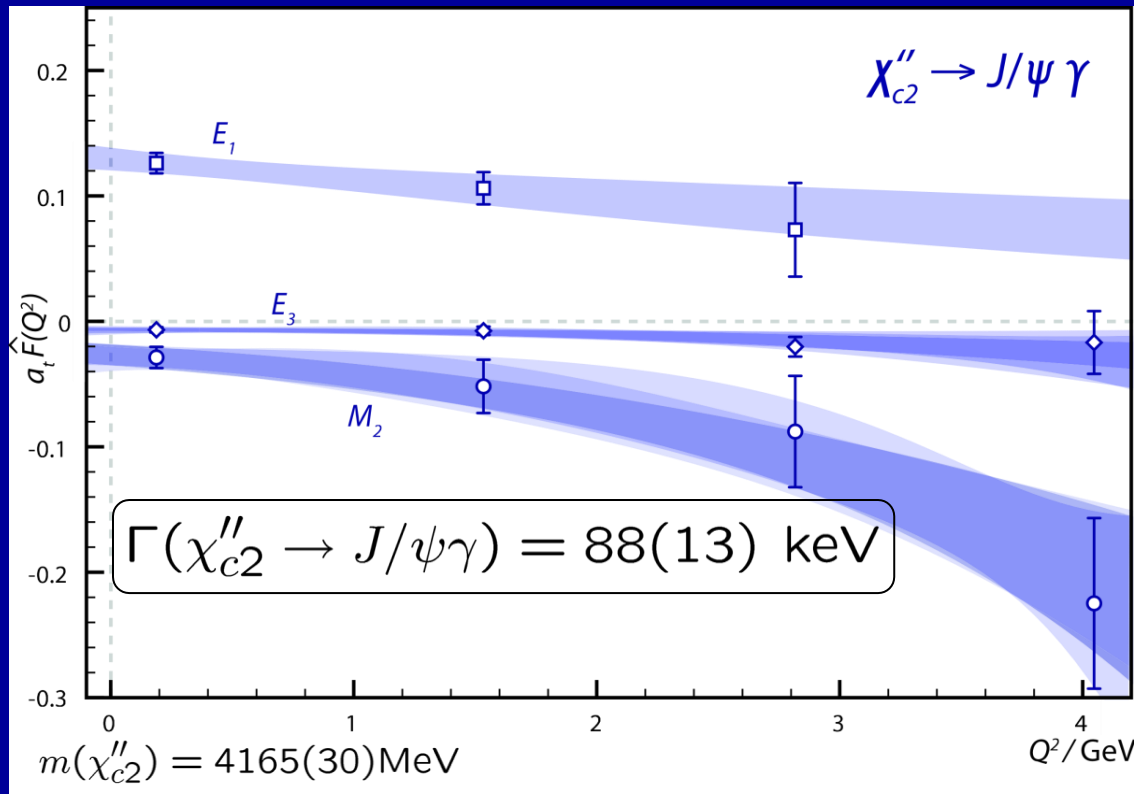


$E_1, M_2, E_3$

Completely different hierarchy!

$|E_3(0)| > |M_2(0)|, |E_1(0)|$

# Tensor $2^{++}$ – Vector $1^-$



$E_1, M_2, E_3$

Quark models ( $2^3P_2$ )  
 $\sim 50 - 80 \text{ keV}$

Reverted to expected hierarchy:

$$|E_1(0)| > |M_2(0)| \gg |E_3(0)|$$

# Tensor $2^{++}$ – Vector $1^-$

Interpretation: **single quark transition** model

In general:  $J_i = J_f \otimes k \quad (k > 0)$   
 $E_1, M_2, E_3 \quad (k = 1, 2, 3)$

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If only a single quark is involved ( ${}^3P_2 \rightarrow {}^3S_1$ ):  
 $j = 3/2 \rightarrow j = 1/2$   
 $k = 1, 2$  only and  $E_3 = 0$   
 $|E_1(0)| > |M_2(0)| \gg |E_3(0)|$

# Tensor $2^{++}$ – Vector $1^{--}$

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$$|E_1(0)| > |M_2(0)| \gg |E_3(0)|$$

If instead tensor is  ${}^3F_2$  ( ${}^3F_2 \rightarrow {}^3S_1$ ):

$$j = 5/2 \rightarrow j = 1/2$$

$$k = 2, 3 \text{ only and } E_1 = 0$$

$$|E_3(0)| > |M_2(0)| \gg |E_1(0)|$$

# Tensor $2^{++}$ – Vector $1^{-}$

Interpretation: **single quark transition** model

In general  $t \rightarrow t' \otimes t''$  ( $t \rightarrow 0$ )

**Interpretation:**

$$\chi_{c2} = 1^3P_2$$

$$\chi'_{c2} = 1^3F_2$$

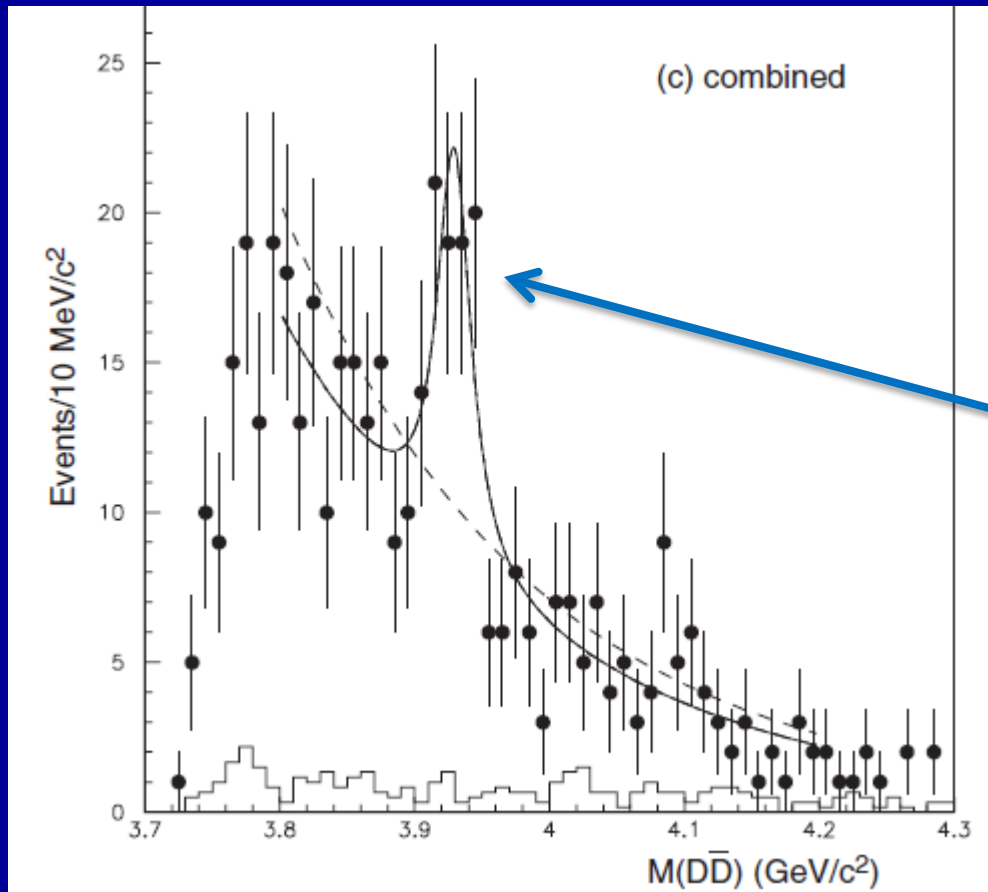
$$\chi''_{c2} = 2^3P_2$$

Supported by spectrum analysis

Squeezed in finite volume?  
Mixing of 1F and 2P?



# Tensor $2^{++}$ – Vector $1^-$



Belle

$\gamma\gamma \rightarrow D\bar{D}$

$\chi'_{c2} \sim 3930$  MeV

Needs lattice calc of  
two-photon coupling  
[extension of Dudek and  
Edwards **PRL 97** 172001  
(2006) ]

Belle [PRL 96 082003 (2006)]

# Vector $1^-$ – Pseudoscalar $0^-$

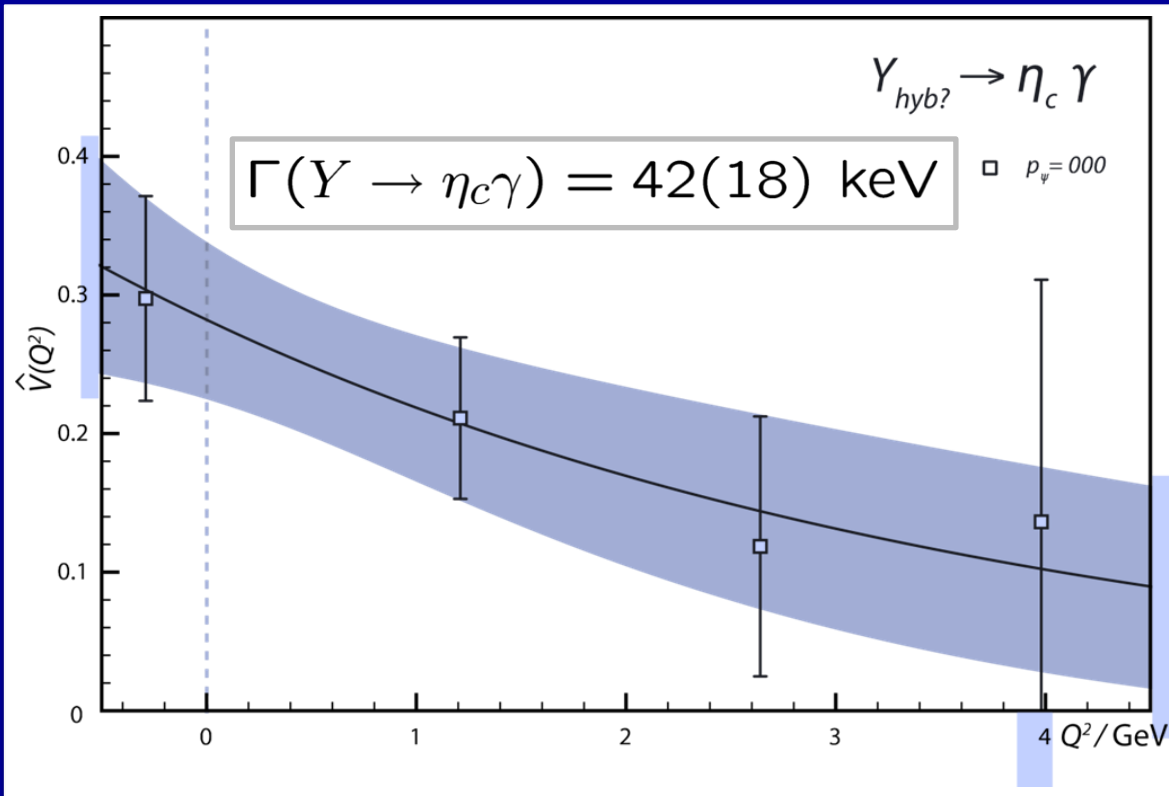
Spectrum results [PR **D77** 034501 (2008) , PR **D78** 094504 (2008) ]:

Level	Mass / MeV	Suggested state	Model assignment
0	3106(2)	$J/\psi$	$1^3S_1$
1	3746(18)	$\psi'(3686)$	$2^3S_1$
2	3846(12)	$\psi_3(3^{--})$	lattice artifact
3	3864(19)	$\psi''(3770)$	$1^3D_1$
4	4283(77)	$\psi('4040')$	$3^3S_1$
5	4400(60)	$Y$	hybrid

Significant overlap with operator  $\sim [D_i, D_j] \sim F$

# Vector $1^-$ – Pseudoscalar $0^+$

Only  $M_1$

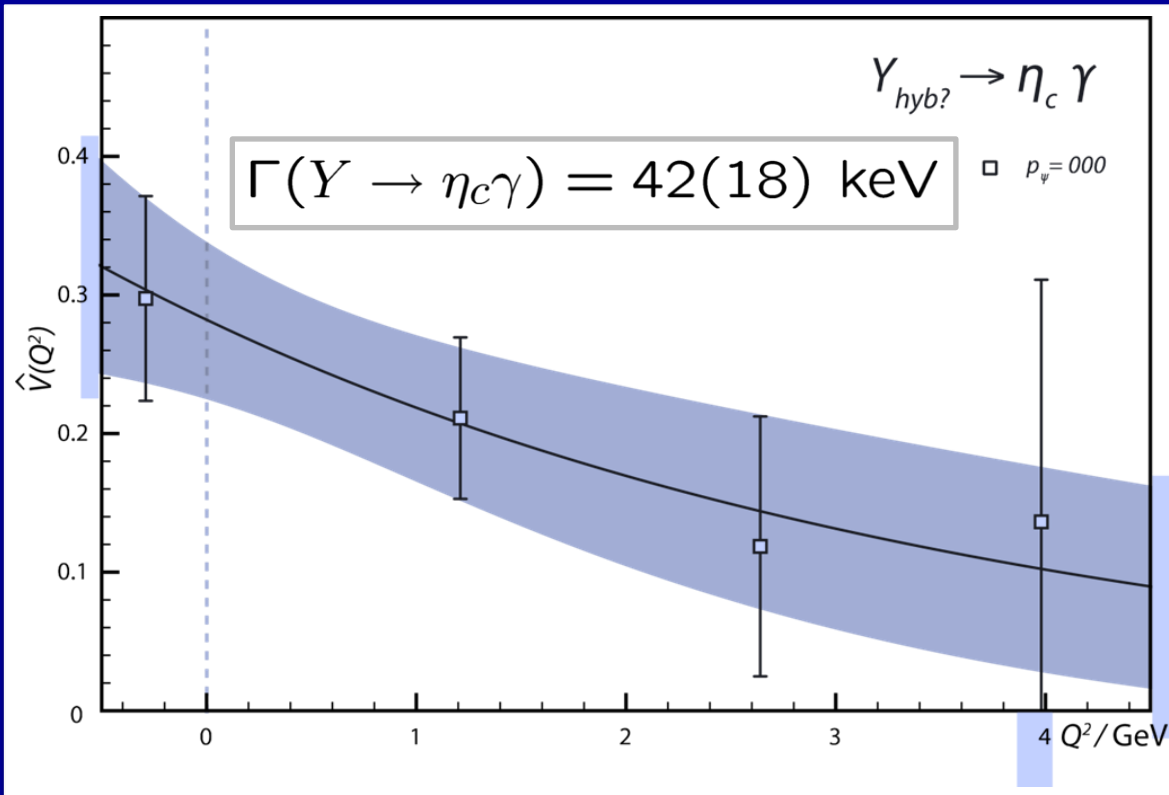


Much larger than other  
 $1^- \rightarrow 0^+$   $M_1$  transitions

$\Gamma(J/\psi \rightarrow \eta_c \gamma) \sim 2 \text{ keV}$

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$$\Gamma(J/\psi \rightarrow \eta_c \gamma) \sim 2 \text{ keV}$$

Spectrum analysis  
suggests a vector hybrid  
(spin-singlet)

Analogous to  $1^-$  hybrid  
to vector trans:  
 $M_1$  with no spin flip

c.f. flux tube model 30 – 60 keV

# Summary and Outlook

## Summary

- **Method successful:** first calc. of excited meson rad. trans. on lattice
- **Hybrid photocoupling is large:**  $\Gamma(\eta_{c1} \rightarrow J/\psi\gamma) \sim 100 \text{ keV}$
- $M_1$  transitions:  $\psi \rightarrow \eta_c\gamma$
- Non-exotic **vector hybrid candidate**  $\Gamma(Y \rightarrow \eta_c\gamma) = 42(18) \text{ keV}$
- $E_1, M_2, E_3$  multipoles;  $2^3P_2, 1^3F_2$  states in  $\chi_{c2} \rightarrow J/\psi\gamma$
- Comparison with models

## Outlook

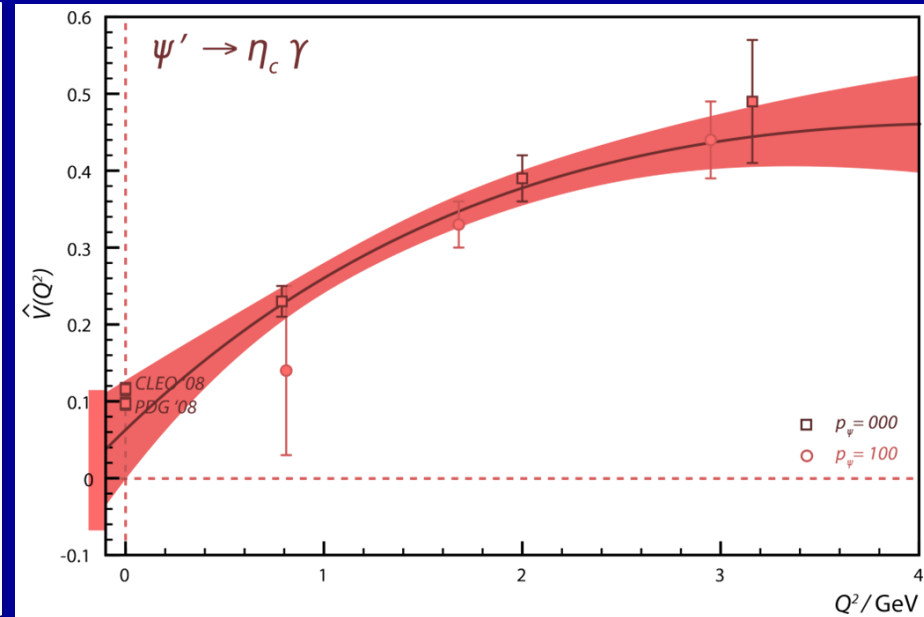
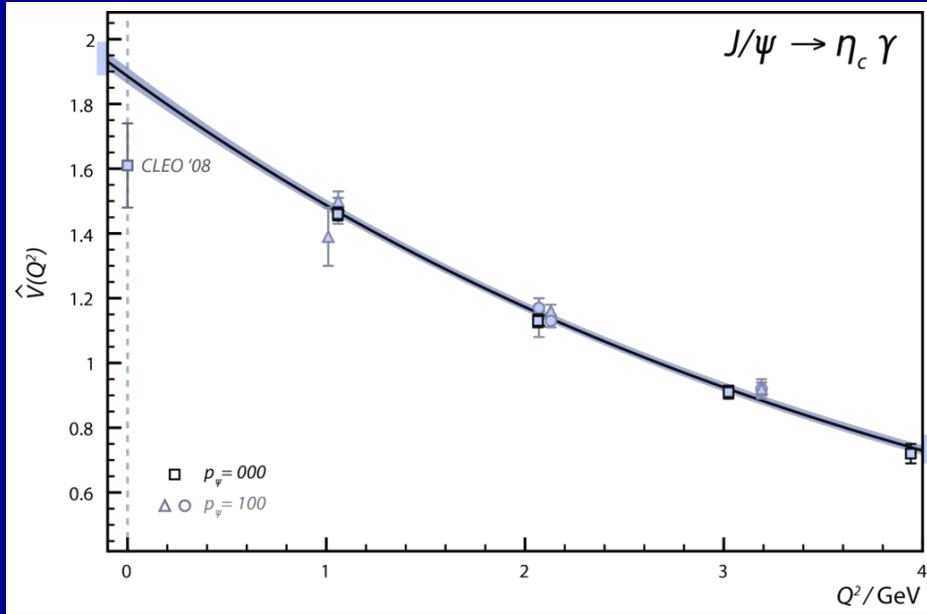
- Systematically improve (unquenched and using 'distillation')
- Apply to lighter mesons (spectrum results) + charmonium (in progress)



# Extra Slides

# Vector $1^-$ – Pseudoscalar $0^-$

Only  $M_1$



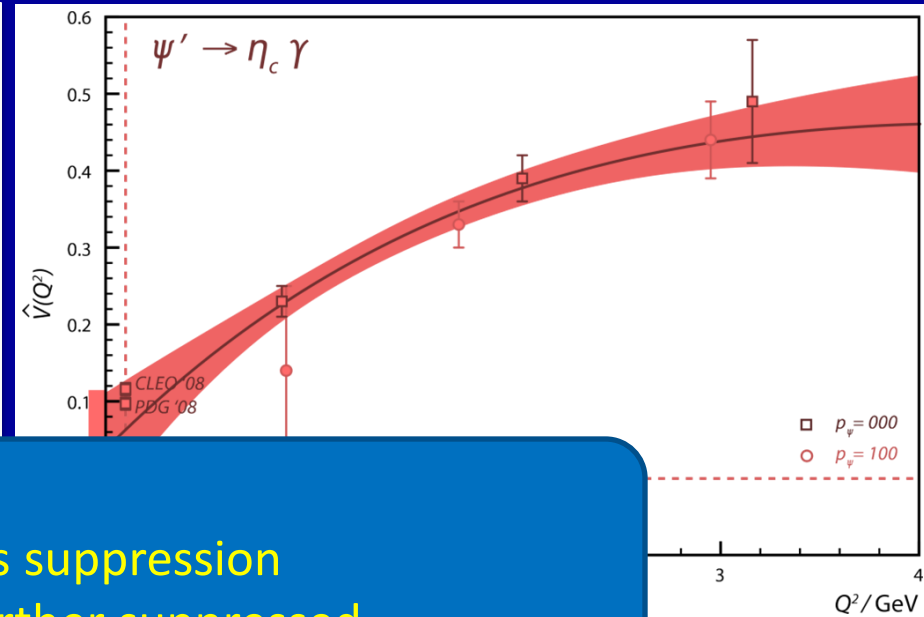
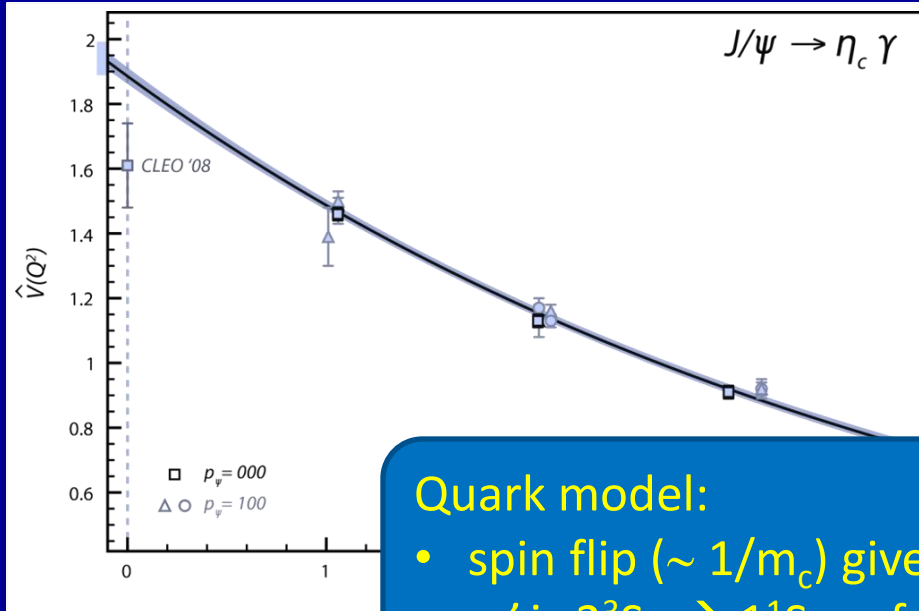
$\Gamma / \text{keV}$	Lattice	Exp.	Barnes, Godfrey, Swanson 'NR'	'GI'	Eichten et. al.
$J/\psi \rightarrow \eta_c \gamma$	2.51(8)	1.85(29) (CLEO-c)	2.9	2.4	1.92
$\psi' \rightarrow \eta_c \gamma$	0.4(8)	0.95(16) (PDG08) 1.37(20) (CLEO-c)	4.6, 9.7	9.6	0.91

[CLEO PRL 102 011801 (2009)]



# Vector $1^-$ – Pseudoscalar $0^-$

Only  $M_1$



Quark model:

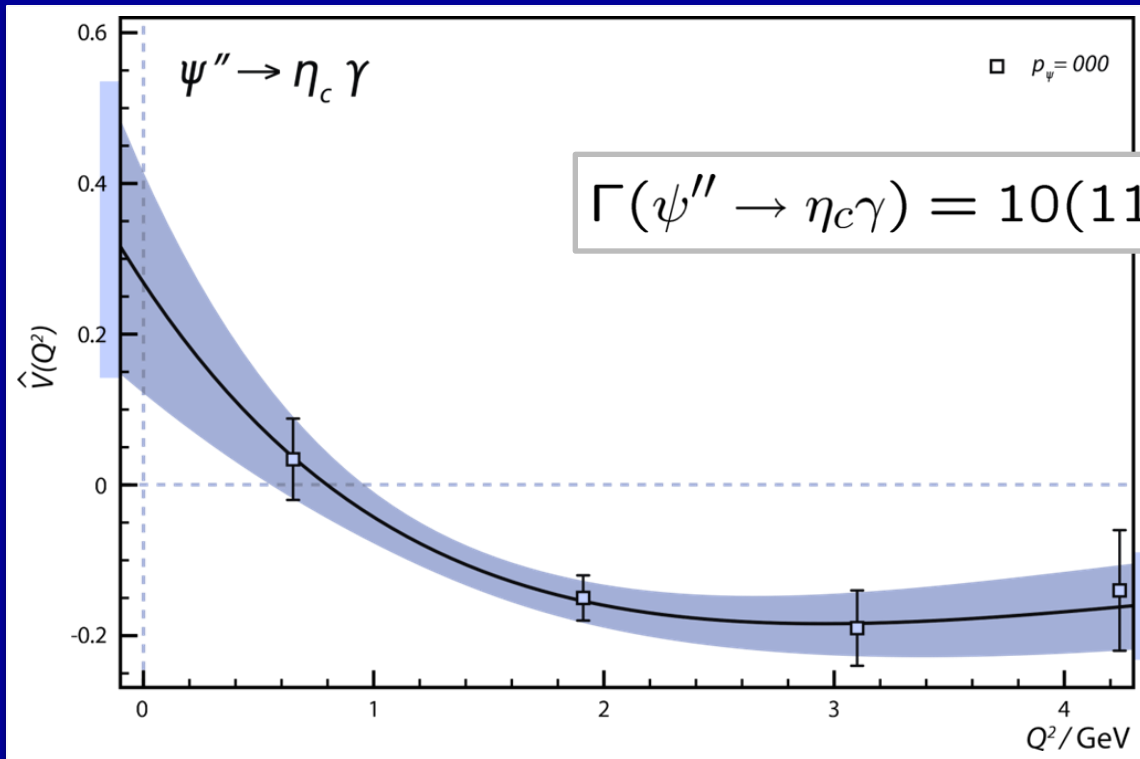
- spin flip ( $\sim 1/m_c$ ) gives suppression
- $\psi'$  is  $2^3S_1 \rightarrow 1^1S_0$  – further suppressed

$\Gamma / \text{keV}$	Lattice	Exp.	Barnes, Godfrey, Swanson 'NR'	'GI'	Eichten et. al.
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[CLEO PRL 102 011801 (2009)]

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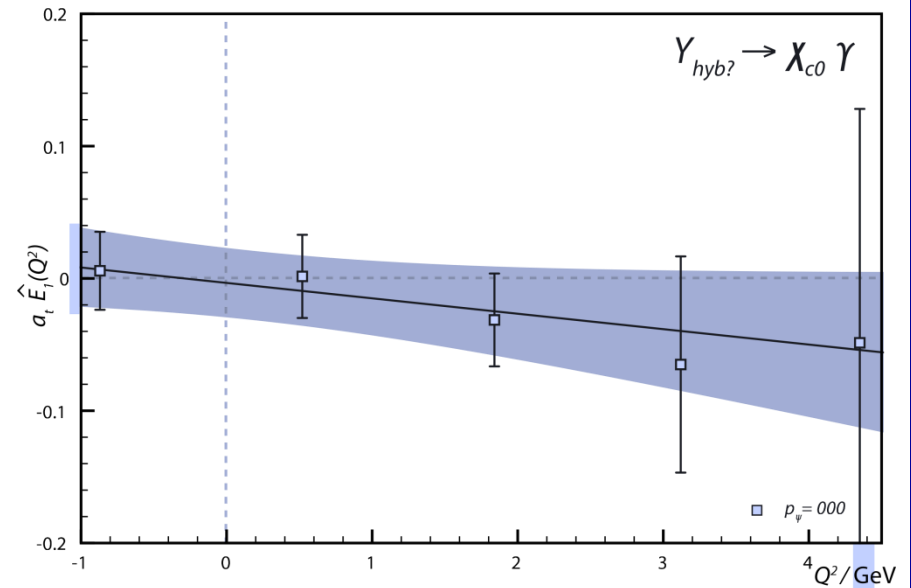
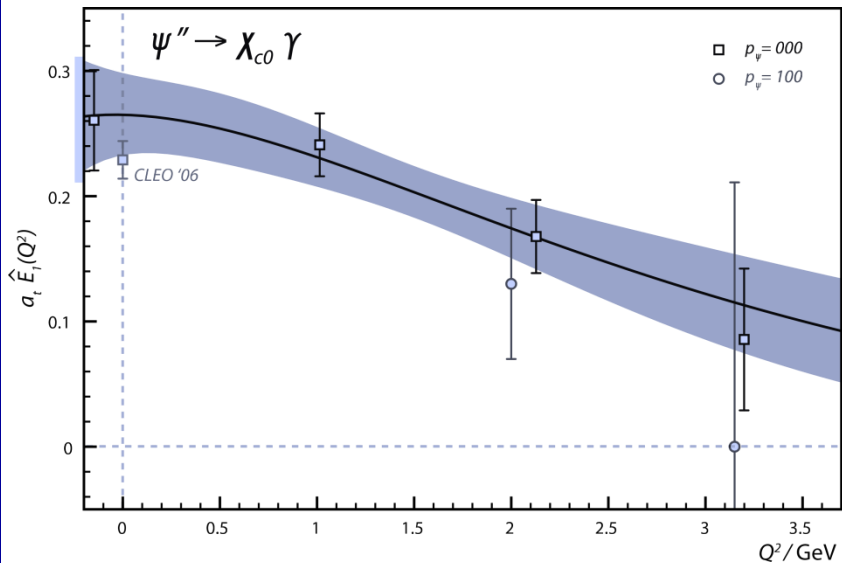
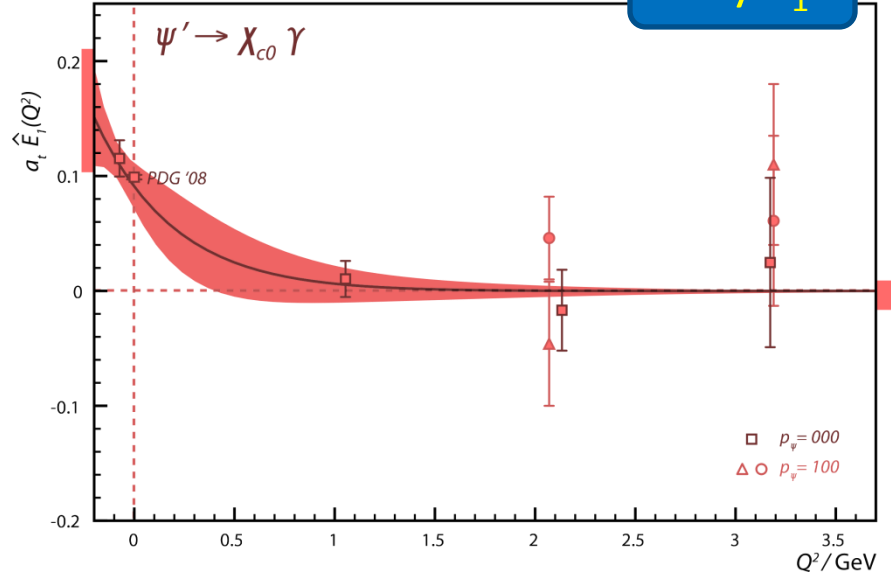
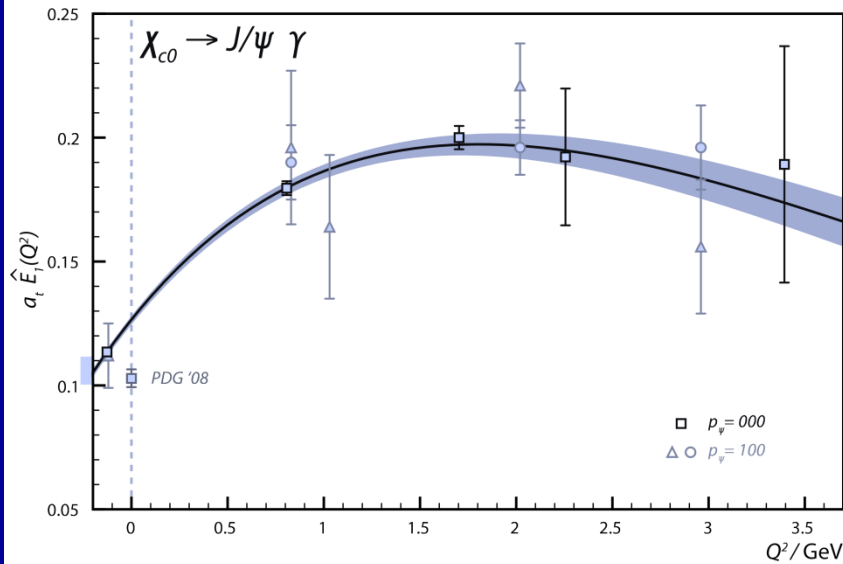
Only  $M_1$



Quark model:  $1^3D_1 \rightarrow 1^1S_0$  has same leading  $Q^2$  behaviour as  $2^3S_1 \rightarrow 1^1S_0$

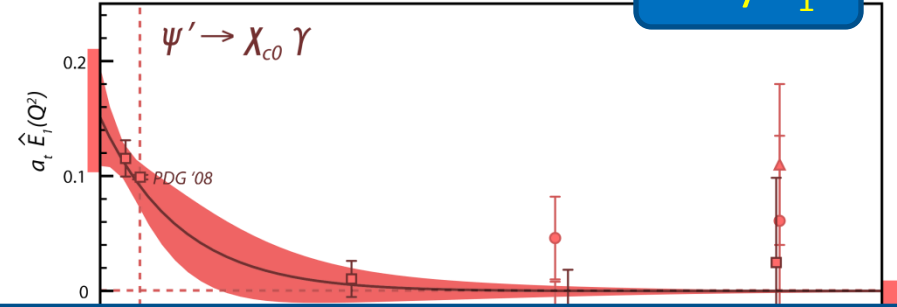
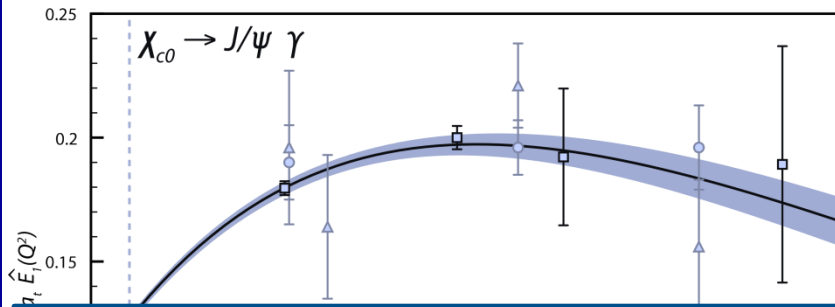
# Scalar $0^{++}$ – Vector $1^{--}$

Only  $E_1$

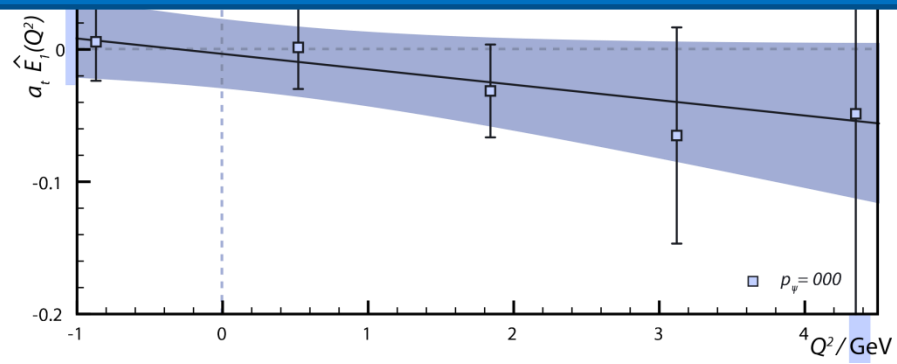
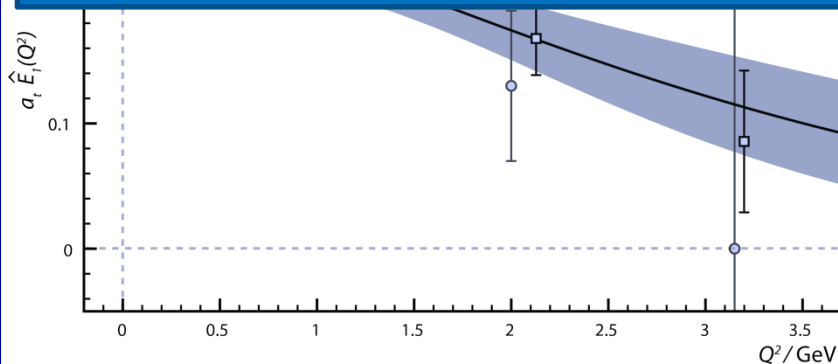


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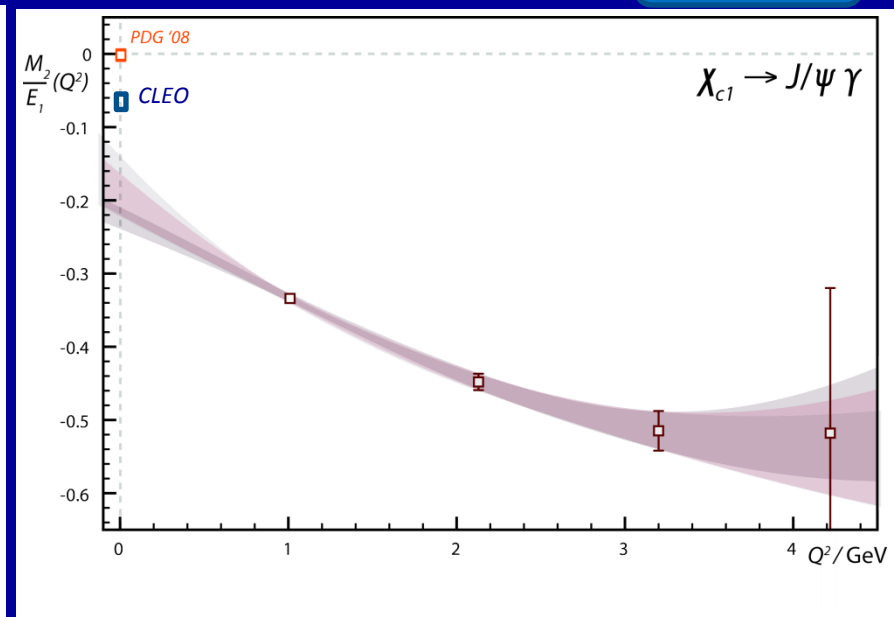
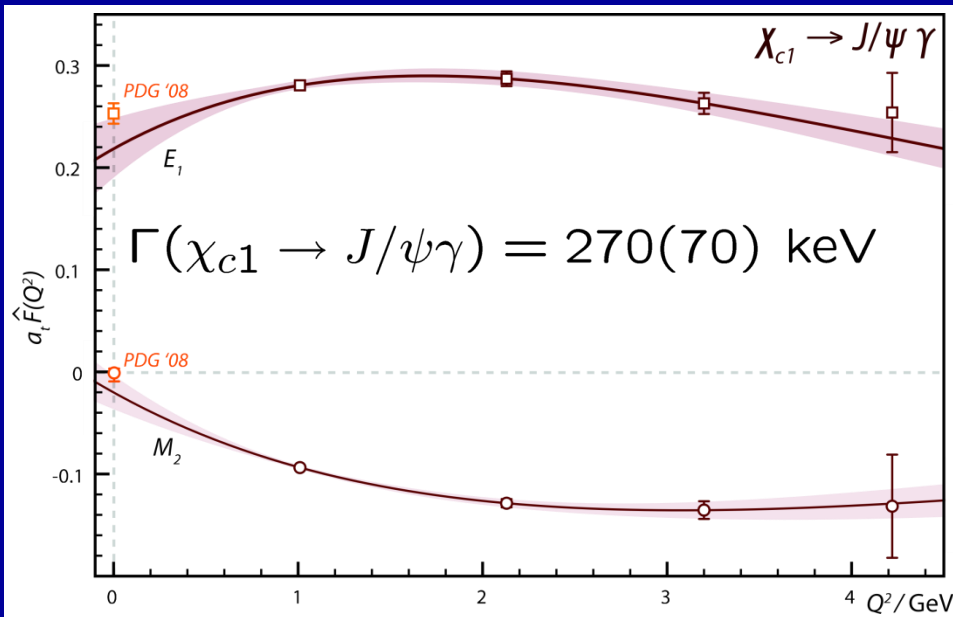


$\Gamma$ / keV	Lattice	Exp. (PDG08)	Barnes, Godfrey, Swanson 'NR'	Barnes, Godfrey, Swanson 'GI'	Eichten et. al.
$\chi_{c0} \rightarrow J/\psi(1^3S_1)\gamma$	199(6)	131(14)	152	114	120, 105
$\psi'(2^3S_1) \rightarrow \chi_{c0}\gamma$	26(11)	30(2)	63	26	46, 38
$\psi''(1^3D_1) \rightarrow \chi_{c0}\gamma$	265(66)	199(26)	403	213	287
$\psi''(3^3S_1) \rightarrow \chi_{c0}\gamma$			0.27	0.63	
$Y \rightarrow \chi_{c0}\gamma$	$\lesssim 20$				



# Axial $1^{++}$ – Vector $1^{-}$

$E_1, M_2$



c.f. PDG08: 320(25) keV

c.f. quark models ( $1^3P_1$ )  $\sim 215 - 314 \text{ keV}$

Expected hierarchy:  $|E_1(0)| > |M_2(0)|$

$$a_2 = M_2 / \sqrt{E_1^2 + M_2^2}$$

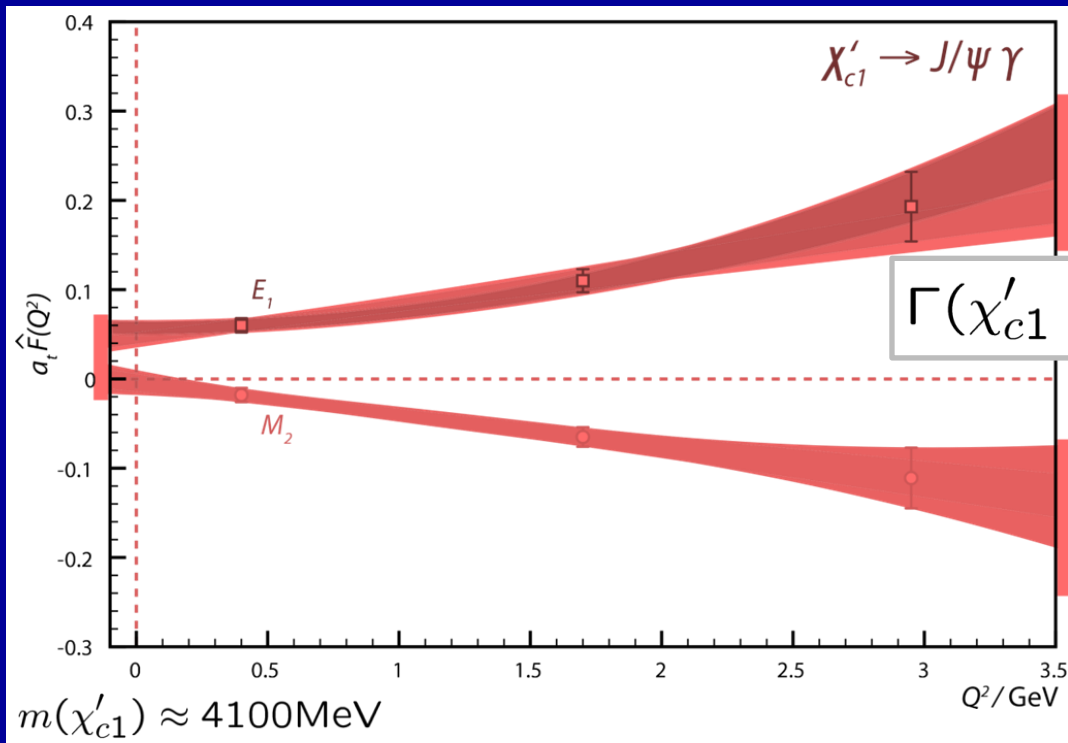
PDG08:  $-0.002^{+0.008}_{-0.017}$

CLEO:  $-0.063(7)$

[CLEO PRD80 112003 (2009)]

# Axial $1^{++}$ – Vector $1^{-}$

$E_1, M_2$



$$\Gamma(\chi_{c1}' \rightarrow J/\psi \gamma) = 21(12) \text{ keV}$$

c.f. quark models ( $2^3P_1$ )  $\sim 14 - 71 \text{ keV}$